

THE IMPORTANCE OF CONCENTRATION EFFECTS AT THE ELECTRODE SURFACE IN ANODIC STRIPPING VOLTAMMETRIC MEASUREMENTS OF COMPLEXATION OF METAL IONS AT NATURAL WATER CONCENTRATIONS

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SUMMARY

The influence of the ligand/metal ion concentration ratio on the shape, peak current and peak potential of curves obtained by anodic stripping voltammetry (a.s.v.) at the hanging mercury drop electrode is described, particularly with respect to the use of a.s.v. for speciation of metal ions at very low concentrations as is often found in natural waters. The lead(II)/triethylenetetramine system is used as a model of a fully labile reversible system. It is shown that the total metal ion concentration at the electrode surface (C_M^0) during the stripping step may be much larger (30–300 times in typical conditions) than that in the bulk solution (C_M), the exact value depending on the deposition time t_d . Consequently, changes in the peak characteristics are observed when the ligand/metal concentration ratio in the bulk of the solution, C_L/C_M , is less than 1000. Semi-empirical equations, experimentally tested, are given, which enable C_M^0/C_M to be estimated for a specified solution and a.s.v. conditions, which correct for the "surface concentration effect" when a.s.v. is used to measure complexation, and which describe the influence of the parameters such as stirring efficiency, radius of the mercury drop and C_L/C_M . The implications of the results are discussed for determinations of total metal ion in complex media, of speciation based on peak-potential shifts or stripping voltammetric curves, and of complexation capacity.

Anodic stripping voltammetry (a.s.v.) is widely used for the determination and speciation of trace metals in natural waters because the techniques can be applied directly, even at the extremely low concentrations found in these samples [1–3]. The a.s.v. procedure includes the deposition step, where the metal ion M is reduced on the electrode surface during time t_d , at a constant deposition potential, E_d , and the stripping step, in which the metal is oxidized back into solution by scanning the potential towards more positive values. The deposition is, of course, a preconcentration step, thus the metal ion concentration at the electrode surface (C_M^0) during the stripping step, is much larger than that in the bulk solution (C_M).

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ligands such as Cl^- , HCO_3^- and SO_4^{2-} , $C_L/C_M \approx 100$ for most of the ligands and C_L/C_M may sometimes be as low as 10. It is then very important to establish whether a ratio between 10 and 100 in the bulk solution is large enough to ensure an excess of ligand at the electrode surface. The aims of this paper are first to evaluate quantitatively the influence of C_L/C_M on certain a.s.v. peak characteristics (current (i_p), potential (E_p), and shape), then to estimate the minimum ratio above which this effect is negligible, and finally to provide a procedure allowing one to correct for this effect when it is not negligible.

In studies of trace metal in natural waters, these aspects are particularly relevant in the following situations: (1) if the total concentration of M is calculated from i_p vs. C_M calibration data, without complete elimination of the ligand in the medium; (2) if the stability constants of labile ML complexes are measured from the shift in E_p with C_L [8]; (3) if complexation of M is evaluated by recording i_p vs. E_d curves (stripping polarography) [9–12]; and (4) if the complexation capacity and complexation constants are calculated from i_p vs. C_M standard addition curves [6, 13–15]. The last three cases will be discussed here.

The above effect was tested by using lead(II) and triethylenetetramine (TETA). In the conditions used, this ligand forms a well defined 1:1 complex [16]. In order to test specifically the surface effects, the bulk composition of the solution was held constant. Hence, in most cases, the increase in C_M^0 was obtained by increasing t_d at constant C_M , and the parameters i_p and E_p were used to follow the resulting effect of surface concentration. This effect was also investigated by using a medium-exchange system, in which the solution in the cell is changed between deposition and stripping.

EXPERIMENTAL

Reagents and apparatus

All reagents were of analytical grade (Merck) except for TETA which was of 99% purity (Fluka). Stock solutions were 0.1 M potassium nitrate acidified to pH 5.5, 1.00×10^{-2} M lead nitrate acidified to pH 3.0, 4.90×10^{-2} M TETA neutralized to pH 8.0 (concentration verified by an acid-base titration), and triethanolamine buffer solution at pH 7.5 with a concentration 10 times that of TETA.

The precision of the pH measurements was ± 0.01 pH unit (Metrohm pH meter E-603). Buffer solutions of pH 4.00 and pH 7.00 (Merck) were used for calibration.

A Tacussel PRG5 polarograph was used for all experiments. The reference electrode, to which all potentials are referred, was Ag/AgCl/sat. KCl/0.10M KNO_3 . The working electrode was a Metrohm (EA-290) hanging mercury drop electrode (HMDE), with a cone-shaped capillary end (the immersed length was 1.7 cm; the outside (o.d.) and inside diameters (i.d.) were 4.55 and 0.1 mm respectively). Because the geometry of the electrochemical cell strongly influences the reproducibility of the deposition step, a special plexi-

glas cell was built (Fig. 2, I) which would ensure reasonably reproducible stirring efficiency. Stirring was provided by a synchronous motor (Metrohm E-504), and a teflon-covered magnetic bar (4 mm diameter), the length of which (1.5 cm) was appropriate to the internal bottom diameter of the cell in order to prevent lateral movement. The bottom of the cell was conical, in order to prevent fallen mercury drops from impeding the rotation of the stirrer. The immersed parts of the platinum counter electrode (1 cm long, 1 mm diameter) and the reference electrode bridge (1 cm long, 3 mm diameter) were positioned about 0.7 cm above the endplane of the HMDE capillary. Purified nitrogen for degassing was introduced by means of a polyethylene capillary tube (1 mm o.d.).

The radius of the mercury drops (r_0), was determined by weighing a fixed number of drops, as a function of the divisions indicated on the HMDE (for 2, 5, 10 divisions, $r_0 = 0.032, 0.044$ and 0.055 ± 0.002 cm, respectively).

The average thickness of the diffusion layer (δ) in the solution during the deposition step was evaluated by recording the current, i_d , in an uncomplexing solution and at a sufficiently negative potential so that $C_{Pb}^0 = 0$, and then by using the equation [17]: $\bar{i}_d = nFS D_{Pb} C_{Pb} / \delta$, where \bar{i}_d is the average value of i_d and is independent of time at sufficiently negative values of E_d (Fig. 3), $n = 2$, $F = 96\,494$ C, S is the surface area of the Hg drop, D_{Pb} is the diffusion coefficient of Pb^{2+} in solution (8.3×10^{-6} cm² s⁻¹) [18], and C_{Pb} is the total concentration of Pb^{2+} in solution (10^{-4} M).

By using a variable-speed motor, it was found that δ increases with $(1/\omega)^{1/2}$, where ω is the rotation speed of the stirrer. With the synchronous

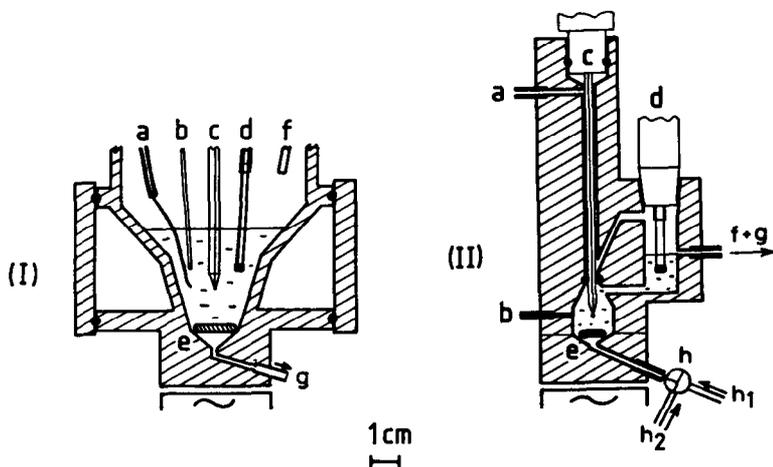


Fig. 2. Plexiglas cells for a.s.v.: (I) under classical conditions; (II) with medium exchange. Parts: (a) nitrogen entrance; (b) Pt electrode; (c) HMDE; (d) reference electrode; (e) stirrer; (f) exit for nitrogen (and solution for cell II); (g) exit of solution (and Hg drops for cell I); (h) three-way stopcock for entry of deposition solution (h_1) and stripping solution (h_2).

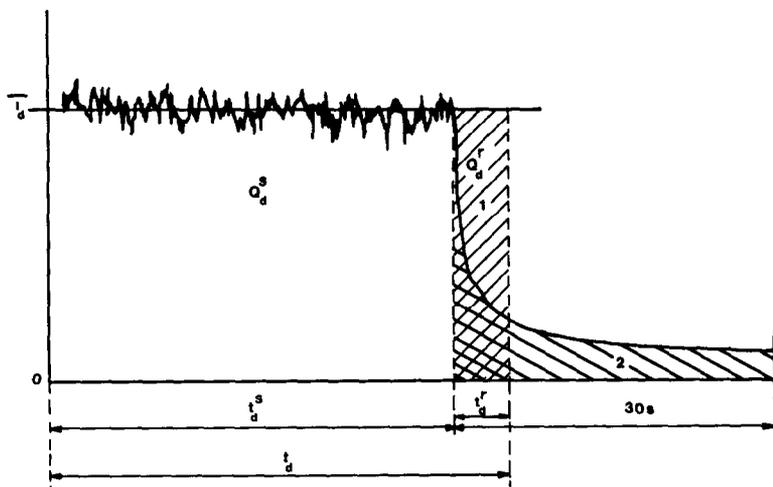


Fig. 3. The reduction current i_d as a function of deposition time t_d . t_d^r is a fictitious time, such that $t_d^r i_d = Q_d^r = \text{area 1} = \text{area 2} = \text{correction term to be applied to } t_d^s \text{ for evaluating the true value of } Q_d = Q_d^s + Q_d^r$.

motor, $\bar{\delta}$ was found to be $(1.2 \pm 0.1) \times 10^{-3}$ cm. This value was measured before each experiment.

The polarographic cell used for medium exchange is shown in Fig. 2(II). Each solution was fully degassed before entering the polarographic cell (inlet h) and a flow of nitrogen was maintained in the cylindrical part around the capillary in order to prevent any penetration of air into the cell. Two reservoirs of degassed solution were connected to inlet h: a "deposition" solution of Pb(II)/TETA (h_1), and a "stripping" solution of 0.10 M potassium nitrate at pH 5.5 (h_2).

Procedure

All measurements were made at $25.0 \pm 0.1^\circ \text{C}$, in 0.10 M potassium nitrate. Experiments in the absence of TETA were done at pH 5.5 in order to prevent the hydrolysis of lead(II). In the presence of TETA, the solutions were buffered with triethanolamine at pH 8.0. It was checked that under these conditions the hydroxide and triethanolamine complexes of lead(II) should be negligible compared to the Pb/TETA complex.

For all the deposition steps, a potential of -0.700 V was applied, while stirring, for a time t_d^s . Then the stirring was stopped, and the solution was allowed to rest for 30 s. The stripping step was then started by scanning the potential in the anodic direction at 20 mV s^{-1} .

As can be seen in Fig. 3, the deposition that continues during the rest period, t_d^r , may be significant, especially when t_d^s is short. The total quantity of electricity corresponding to the deposition is given by $Q_d = Q_d^s + Q_d^r$,

where Q_d^s and Q_d^r are the quantities of electricity corresponding to the stirring and rest periods, respectively. Dividing through by \bar{i}_d , the constant average deposition current, gives $Q_d/\bar{i}_d = Q_d^s/\bar{i}_d + Q_d^r/\bar{i}_d$ or $t_d = t_d^s + t_d^r$. Here t_d is the "corrected" deposition time proportional to the total amount of lead deposited in the drop, and t_d^s is the correction term. By recording curves similar to that in Fig. 3, for $C_L = 0$ and various concentrations of lead ions ($C_{Pb} = 10^{-4}$ M, 10^{-6} M, and $\approx 10^{-8}$ M, corresponding to the concentration of lead(II) contaminations), it was observed that t_d^s is independent of C_{Pb} and depends only on the mode of stirring. In the present conditions, it was found that $t_d^s = 5$ s. This correction was applied in all cases and only t_d values are reported below. The same type of correction was applied in the case of medium exchange, but t_d^s was found to be larger (25 s).

For the determination of the stability constants of the Pb/TETA complex, all E_p values (for cathodic sweep or a.s.v. techniques) in the presence of TETA, and the corresponding values in its absence, to which they are referred, were measured within the same set of experiments. All errors reported below are standard deviations, σ , defined as $\sigma = [\sum_i (x_i - \bar{x})^2 / (N - 1)]^{1/2}$, where N is the number of determinations.

Electrochemical characteristics of the lead(II)/TETA system

The reported values of the acid-base constants of TETA and of its stability constant with lead(II), are given in Table 1. Cathodic sweep voltammetry was used to test the electrochemical behaviour of the Pb(II)/TETA system with $C_{Pb} = 10^{-5}$ M, $C_L = 9.1 \times 10^{-3}$ M and pH 8.02. According to the usual criteria [19], ($E_p - E_{p/2} = 28$ mV, and independence of E_p , $i_p/v^{1/2}$ and $E_p - E_{p/2}$ from the scan rate v , $E_{p/2}$ being the potential corresponding to $i_{p/2}$) it behaves as a labile reversible system, in the range $10 < v < 40$ mV s⁻¹.

For both a.s.v. and cathodic sweep voltammetry, the ratio of the peak currents, in the absence (pH 5.5) and presence (pH 8) of TETA, i_p^{nc}/i_p^c ,

TABLE 1

Characteristics of the lead(II)/TETA system^a

Reaction	T (°C)	Constant	Logarithm of stability constant	Ref.
$L + H \rightleftharpoons LH$	25	β_1^H	9.74	16
$L + 2H \rightleftharpoons LH_2$	25	β_2^H	18.82	16
$L + Pb \rightleftharpoons PbL$	20	β	10.4	16
$L + Pb \rightleftharpoons PbL$	25	β	10.18 ± 0.03^b	This work

^aTETA = L = NH₂-CH₂-CH₂-NH-CH₂-CH₂-NH-CH₂-CH₂-NH₂. All data are for an ionic strength of 0.1 M. ^bAverage of four determinations by differential-pulse cathodic sweep voltammetry with $C_{Pb} = 2 \times 10^{-5}$ M and C_L varying between 0.98×10^{-3} and 1.95×10^{-2} M.

were found to be 1.00 ± 0.05 , and were independent of the values of C_{Pb} and C_L , provided that TETA was in large excess (>30 for cathodic sweep voltammetry and >1000 for a.s.v.). This result implies that $D_{PbL} = D_{Pb}$. Consequently, the shifts in E_p in cathodic sweep voltammetry and in differential-pulse cathodic sweep voltammetry were used, as a further test of the electrochemical behaviour of the system, to compute the stability constant of the Pb(II)/TETA complex, according to Eqns. 2 and 3 (see below). The value shown in Table 1 is slightly smaller than the literature value, which may be due to the difference in temperature. In any case, considering the small amount of data available in the literature, the correspondence of the data in Table 1 was considered as satisfactory. In particular, the good agreement between the cathodic sweep voltammetry and the differential-pulse results, both for i_p^{nc}/i_p^c and $E_p^c - E_p^{nc}$, is a further confirmation of the lability of the complexation reaction. The suffixes nc and c stand for noncomplexing ($C_L = 0$) and complexing ($C_L \neq 0$) media respectively.

THEORY, DEFINITIONS AND QUALITATIVE DESCRIPTION

Discriminating between complexation at the electrode surface and in the bulk solution

For the lead(II)/TETA system described above, the degree of complexation of lead(II) in the bulk solution is given [20] by

$$\alpha = C_{Pb}/[Pb^{2+}] = 1 + \beta[L] = 1 + \beta(C_L - [PbL])/ \alpha_H \quad (1)$$

where $\alpha_H = 1 + \beta_1^H[H] + \beta_2^H[H]^2$; C_{Pb} and C_L are the total concentrations of lead(II) and TETA, and $[Pb^{2+}]$, $[L]$ and $[H]$ are the free metal ion, ligand and proton concentrations. When a.s.v. is used to quantify lead(II) in this system, the resulting curve depends on the degree of complexation of lead(II) at the electrode surface, both during the deposition step (α_d^0), and the stripping step (α_s^0). Therefore, it is easy to compute the value of α from the characteristic parameters of the peak, E_p and i_p , only when $\alpha = \alpha_d^0 = \alpha_s^0$. As will be shown below, this is valid when $C_L \gg C_M$. Then Eqn. 1 reduces to

$$\alpha = 1 + \beta C_L / \alpha_H \quad (2)$$

In such a case, for a labile, reversibly reduced system like Pb(II)/TETA, the DeFord and Hume method [21, 22], initially developed for d.c. polarography, can be applied to the a.s.v. technique

$$\ln \alpha = (nF/RT)(E_p^{nc} - E_p^c) + \ln (i_p^{nc}/i_p^c) \quad (3)$$

In the particular case of the lead(II)/TETA system, it was observed that $i_p^{nc} = i_p^c$ so that the second term in Eqn. 3 vanishes. By combining Eqns. 2 and 3, β can easily be calculated, when α is a measured parameter.

Condition for $\alpha = \alpha_d^0$ during the deposition step

The reduction of complexes without an excess of ligand in d.c. polarography has been discussed in detail [22]. For the reduction of metal complexes

for reversible systems in the absence of ligands for mercury film and hanging drop electrodes. For the HMDE, when E_d is chosen on the plateau of the reduction wave, and the scan rate of the stripping step is large enough ($>5 \text{ mV s}^{-1}$) so that the bulk concentration of M^0 in the mercury drop is not changed during the stripping, it was shown that at the peak potential, E_p

$$C_M^0/C_M = (1/0.44)(D_{Ox}D_R)^{1/2} t_d/r_0\delta \quad (4)$$

where D_{Ox} and D_R are the diffusion coefficients of the metal ion in solution and of the metal in mercury, respectively. It should be noted that C_M^0/C_M is independent of v , in contrast to i_p which increases with $v^{1/2}$ [18]. From Eqn. 4, using typical parameter values and with $t_d = 3 \text{ min}$, it is possible to estimate that $C_M^0/C_M \approx 50$, i.e., surface concentrations much larger than bulk ones may be reached during the stripping step. Figure 4(b) shows the concentration gradients of M and L at E_p , when the stripping occurs in the presence of an "excess" of complexing agent: (e.g., $C_L/C_M = 30$): for a short t_d , $C_L^0 \approx C_L$ is still much larger than C_M^0 and $\alpha_s^0 = \alpha$. However, for large t_d , C_M^0 may become larger than C_L^0 so that $\alpha_s^0 \ll \alpha$. From the above estimate of C_M^0/C_M , it may be expected that $C_L/C_M > 1000$ is needed in the bulk solution, in order to ensure that $C_L^0/C_M^0 > 20$ at the electrode surface (for $t_d = 3 \text{ min}$).

A qualitative description of the shape of the a.s.v. curves can be understood from Figs. 4(b, I) and 5. Figure 4 (b, I) shows the concentration gradients of M at different potentials during the stripping step for a given value of t_d . The surface concentration of M increases in the potential range corresponding to the rising portion of the a.s.v. curve, and remains constant from then on (cf. [5]). When t_d is small enough, $C_L^0/C_M^0 \approx C_L/C_M$ for all potentials, and a peak with a normal shape is obtained. For intermediate values of t_d , $C_L^0 > C_M^0$ for all potentials, but C_L^0/C_M^0 is much smaller at E_p than at the foot of the peak (see Fig. 4b, I). As a consequence, α_s^0 decreases all along the peak, and the resulting peak is broadened (Fig. 5, curve 1).

Finally, for still larger values of t_d , L becomes fully saturated at the electrode surface ($C_L^0 \approx C_M^0$) at a potential, $E_{sat} < E_p$. For $E > E_{sat}$, $\alpha_s^0 \approx 1$, thus the oxidation is greatly decreased, until the potential reaches a value corresponding to that for the oxidation in an uncomplexing medium. In such cases, a shoulder appears in the a.s.v. curves (Fig. 5, curves 2 and 3) for systems such that $E_p^{nc} - E_p^c(0)$ is large enough ($E_p^c(0)$ is the value of E_p^c at $t_d = 0$, see Fig. 6). The height of the shoulder, i_{sh} , depends mainly on C_L , and does not vary when t_d increases (Fig. 6). Hence, for extremely large values of t_d , the largest part of M^0 is reoxidized as free M^{2+} , giving rise to the corresponding a.s.v. peak, and the relative contribution of the shoulder to the total current becomes increasingly smaller, although its absolute value remains constant.

The influence of t_d on the a.s.v. peak parameters, E_p^c , i_p^c and i_{sh} , as described above, is shown in Fig. 6.

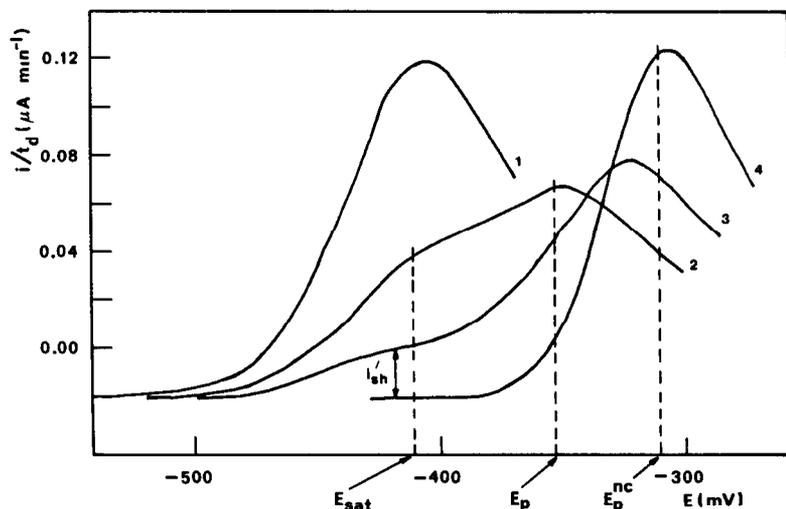


Fig. 5. "Normalized" anodic stripping voltammograms (i/t_d vs. potential) of lead ion ($C_M = 1.00 \times 10^{-6}$ M): (1–3) in presence of TETA ($C_L = 0.98 \times 10^{-4}$ M) for $t_d = 1, 7$ and 25 min, respectively; (4) in the absence of TETA. Other experimental conditions: 10^{-3} M triethanolamine, pH 8.02, $\mu = 0.10$ M (KNO_3); $T = 25.0^\circ\text{C}$; $r_0 = 0.032$ cm; $v = 20$ mV s^{-1} ; $\delta = 1.2 \times 10^{-3}$ cm. E_{sat} and E_p are given for curve 2 (see text) $i'_{\text{sh}} = i_{\text{sh}}/t_d$.

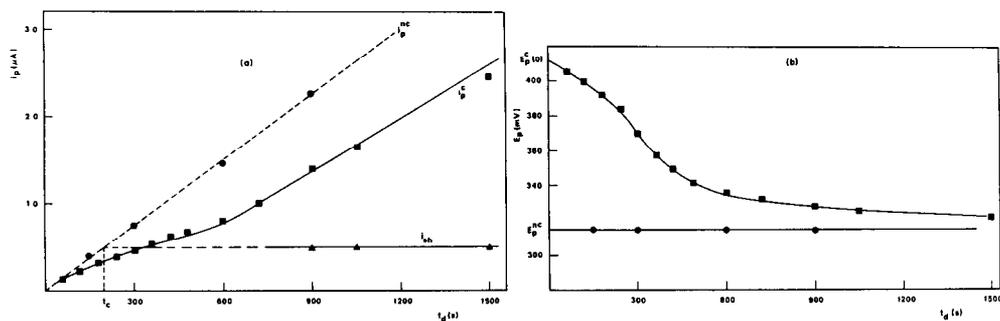


Fig. 6. Changes in the characteristic parameters of a.s.v. peaks (E_p , i_p) as a function of t_d : (●) points corresponding to the a.s.v. peak of the metal in the absence of the ligand; (■) points corresponding to the a.s.v. peak of the metal in the presence of the ligand; (▲) points corresponding to the a.s.v. peak shoulder height i_{sh} . Experimental conditions as in Fig. 5.

Definition of critical deposition time

From the above description, it is clear that the ratio C_M^0/C_M is a critical factor for correct interpretation of the a.s.v. parameters (E_p or i_p) obtained in any complexing medium. At present, there is no theoretical relationship which allows one to predict the values of i_p , E_p , or C_M^0/C_M , in conditions where there is no excess of ligand at the electrode surface during the stripping step.

The usefulness of Eqn. 4 for this purpose was therefore tested in the following part of this work. Although it is rigorously valid only in the absence of ligands, it may be expected that the influence of each of its parameters should be similar in complexing media and that it could be used as a semi-empirical equation for analytical prediction. In this respect, the lead(II)/TETA system is very suitable, because it was shown that the diffusion coefficients of the free and complexed forms of lead(II) are equal.

In order to test Eqn. 4 rigorously, it would be necessary to measure C_M^0 as a function of t_d . This is not possible, and so use was made of the fact that a shoulder appears in the a.s.v. curves when $C_M^0 \approx C_L$ and remains constant for $C_M^0 > C_L$, even for a very large t_d when $C_M^0 \gg C_L$ (Figs. 5 and 6). The height of this shoulder (i_{sh}) was then used to define a critical deposition time, t_c (as shown in Fig. 6), corresponding to the minimum value of t_d for which TETA is considered to be saturated at the electrode surface. Then one has $C_L^0 = C_L = C_M^0$, and

$$t_c = K(r_0\delta/D_{Ox}D_R)(C_L/C_M) \quad (5)$$

The experimental measurement of t_c was chosen so that the condition $C_L^0 = C_L = C_M^0$ was reasonably fulfilled. Even though this may not be rigorously true, t_c can be considered as a useful semiempirical parameter for testing the influence of the factors in Eqn. 5. Because Eqn. 4 was not deduced for the presence of ligands, the numerical constant K could be different from 0.44, but its order of magnitude should be similar. The influence of the selected parameter on the interpretation of the results was tested by defining another critical deposition time, $t_c(E)$, from the E_p^c vs. t_d curve (Fig. 6). This curve may be considered as a compleximetric titration curve of L by M inside the diffusion layer, the potential jump corresponding to the end-point of the titration. Because the break in this curve is always more clearly measurable than the inflection point, it was used to define $t_c(E)$ (Fig. 6). In all cases tested, the effects of the parameters of Eqn. 4 on t_c and $t_c(E)$ were similar. However, the measurement of $t_c(E)$ was less precise and only the results obtained with t_c are reported here.

RESULTS

Influence of r_0 , $\bar{\delta}$, and C_L/C_M

The changes in t_c with r_0 , $\bar{\delta}$, and C_L/C_M are given in Table 2 and Fig. 7. As predicted by Eqn. 5, a linear relationship is observed between t_c and $r_0\bar{\delta}(C_L/C_M)$. From the slope of the line $(6.0 \pm 0.4) \times 10^4$ s cm⁻², with the values $D_{Ox} = 8.3 \times 10^{-6}$ cm² s⁻¹ and $D_R = 1.24 \times 10^{-5}$ cm² s⁻¹ [23], it was calculated that $K = 0.61 \pm 0.04$. From the discussion of the critical deposition time this value of K is in reasonably good agreement with the theoretical value of 0.44 in Eqn. 4. This result confirms the quantitative influence of r_0 , $\bar{\delta}$, and C_L/C_M on the a.s.v. peak parameters, and shows that Eqn. 4 provides an estimate of C_M^0 , even in the presence of a ligand.

TABLE 2

Values of t_c found for different values of the experimental parameters r_0 , $\bar{\delta}$ and C_L/C_M (errors were $r_0 = \pm 0.002$; $\bar{\delta} = \pm 10\%$; $t_c = \pm 12$ s)

r_0 (cm)	$\bar{\delta}$ (cm)	C_L/C_M	t_c (min)
0.055	1.2×10^{-3}	98	6.2
0.044	1.2×10^{-3}	136	7.1
0.044	1.2×10^{-3}	98	5.2
0.044	1.2×10^{-3}	65	3.4
0.044	1.7×10^{-3}	98	7.6
0.044	2.3×10^{-3}	98	11.0
0.032	1.2×10^{-3}	98	3.4

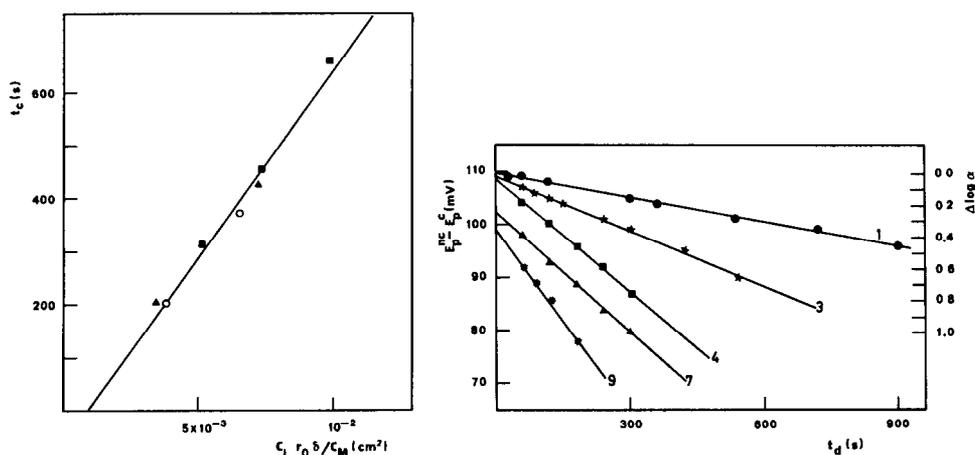


Fig. 7. Influence of the experimental parameters $\bar{\delta}$, r_0 , C_L/C_M , given in Table 2, on the critical deposition time t_c . (The point at $t_c = 11$ min is less reliable, because a very long electrolysis time (45 min) was used for measurement of i_{sh} . (\circ , \square , \blacktriangle) Variations of r_0 , $\bar{\delta}$, and C_L/C_M , respectively (see Table 2).

Fig. 8. Change of $E_p^{\text{nc}} - E_p^{\text{c}}$ and $\Delta \log(\alpha)$ with t_d , for $t_d < t_c$. Plots correspond to the experimental conditions presented in Table 3 with the same number: (\bullet) 1; (\star) 3; (\blacksquare) 4; (\blacktriangle) 7; (\ast) 9. The axis $\Delta \log \alpha$ relates to line 3 only.

The E_p^{c} vs. t_d curves, and measurements of β

Figure 8 shows the change in E_p^{c} with t_d , expressed as $\Delta E_p = E_p^{\text{nc}} - E_p^{\text{c}}$, for $t_d < t_c$. Linear relationships were always obtained within experimental error. This property makes it possible to evaluate β by extrapolating ΔE_p to $t_d = 0$ [$\Delta E_p(t = 0)$], i.e., to conditions where $\alpha_s^0 = \alpha$. The $\Delta E_p(t = 0)$ values were then used with Eqns. 2 and 3 to compute β . The values obtained in this way for various values of r_0 , $\bar{\delta}$, C_L/C_M and pH, are reported in Table 3. All these values are statistically identical, within experimental error, and the average

TABLE 3

Stability constant ($\log \beta$) evaluated from the intercept of the linear relations (ΔE_p , t_d) with slopes α , for different values of the experimental parameters r_0 , $\bar{\delta}$, C_L , C_L/C_M and pH^a

No.	$\bar{\delta}$ (cm)	C_L (10^{-4} M)	C_L/C_M	pH	Slope ($\mu V s^{-1}$)	$\log \beta$
1	1.2×10^{-3}	4.90	1225	7.80	15.3 ± 0.3	10.28 ± 0.01
2	1.2×10^{-3}	4.90	612	7.81	23.2 ± 0.3	10.27 ± 0.01
3	1.2×10^{-3}	4.90	306	7.81	34.8 ± 0.7	10.25 ± 0.01
4	1.2×10^{-3}	4.90	136	7.81	72 ± 2	10.25 ± 0.01
5	1.2×10^{-3}	4.90	65	7.80	137 ± 8	10.21 ± 0.03
6	1.2×10^{-3}	0.98	98	8.05	77 ± 8	10.23 ± 0.06
7	1.7×10^{-3}	0.98	98	8.05	76 ± 2	10.25 ± 0.01
8	2.3×10^{-3}	0.98	98	8.08	43 ± 2	10.27 ± 0.02
9	1.2×10^{-3}	0.98	98	8.02	117 ± 7	10.20 ± 0.03
10	1.2×10^{-3}	0.98	98	8.06	88 ± 5	10.27 ± 0.03

^aScan rate 20 mV s^{-1} , $r_0 = 0.044 \text{ cm}$ for all except nos. 9 and 10, for which $r_0 = 0.032$ and 0.055 cm , respectively.

value, $\log \beta = 10.25 \pm 0.03$, is in agreement with that found by differential-pulse cathodic sweep voltammetry (Table 1).

It must be pointed out that in the present a.s.v. determination of $\log \beta$, the chief sources of error are those incurred from the measurement of $\Delta E_p(t = 0)$ ($\pm 2 \text{ mV}$), and pH (± 0.01). These errors produce errors of 0.07 and 0.02 log unit in β , respectively.

Figure 8 shows that not only does E_p^c increase linearly with t_d (for $t_d < t_c$), but that the slopes of the lines depend on the same parameters as t_c . In this time range, TETA never reaches saturation at the surface, but α_s^0 decreases gradually when t_d increases. On the basis of this property, it is possible to derive a semiempirical equation (see Appendix) for the change in ΔE_p with t_d

$$\Delta E_p = \Delta E_p(t = 0) - (1/K')(RT/nF)(D_{Ox}D_R)^{1/2}(C_M/C_L)t_d/r_0\bar{\delta} \quad (6)$$

with $\Delta E_p(t = 0) = (RT/nF) \ln(\beta C_L/\alpha_H)$; K' is an empirical constant, and should have a value close to 0.44 (Eqn. 4).

The slopes of the straight lines ΔE_p vs. t_d , are given in Table 3, and are plotted as a function of $C_M/(C_L r_0 \bar{\delta})$ in Fig. 9. A linear relationship is obtained with a slope of $(5.7 \pm 1.8) \times 10^{-4} \text{ mV cm}^2 \text{ s}^{-1}$. Then by using the above values of D_{Ox} and D_R , the value of K' obtained is 0.50 ± 0.16 .

Effects of medium exchange on the a.s.v. curves

The medium-exchange system was used to verify that changes in the a.s.v. peak parameters with t_d , were due to phenomena occurring during the stripping step, and not during the deposition step.

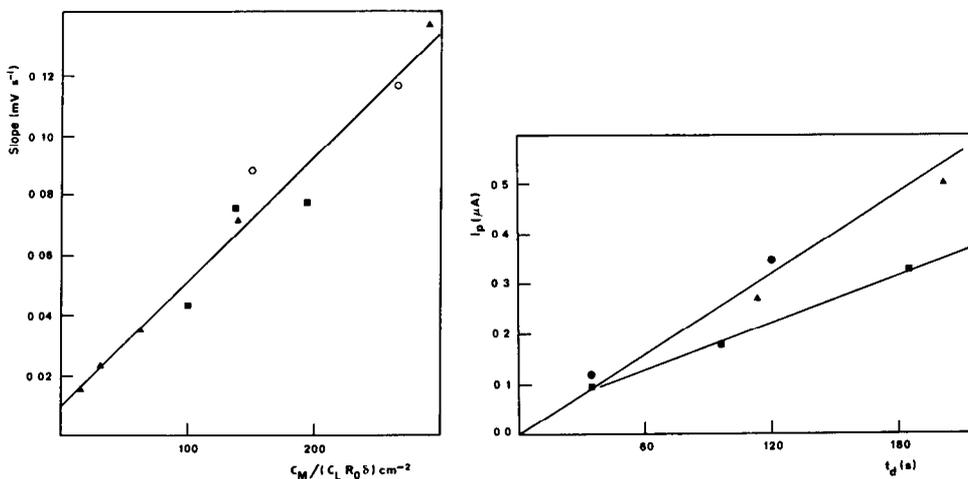


Fig. 9. Influence of the experimental conditions $\bar{\delta}$, r_0 and C_L/C_M on the slopes of the ΔE_p vs. t_d data. Experimental conditions are given in Table 3. The symbols (\circ , \blacksquare , \blacktriangle) are the same as in Fig. 7.

Fig. 10. Peak current i_p vs. t_d of the a.s.v. curves: (\bullet) in the absence of TETA; (\blacksquare) in the presence of TETA, without changing the solution; (\blacktriangle) in the presence of TETA, by changing the solution. Each point is the average value of 3–7 determinations. In all cases, experimental errors on i_p varied between 1 and 5% and $C_{Pb} = 1.0 \times 10^{-6}$ M. For (\blacksquare) and (\blacktriangle), $C_L = 4.6 \times 10^{-5}$ M.

Figure 10 shows i_p vs. t_d relationships in the absence of TETA, and in the presence of TETA without and with medium exchange after the deposition step. It is clear that, without medium exchange, an i_p vs. t_d curve similar to that in Fig. 6 is obtained, whereas, with medium exchange, the points fall on the same straight line as that obtained in the absence of TETA. Not only does this result confirm that the observations described above are due to the stripping step, but also that the diffusion rate of Pb^{2+} and PbL during the deposition step are identical.

DISCUSSION

The results described above show the necessity of considering the concentrations of reagents at the electrode surface when a.s.v. techniques are applied in media containing complexes. The dependences of i_p and E_p on t_d show that the influences of the important factors on C_M^0 are adequately predicted by Eqn. 4 and that this equation can even give a reasonably good estimate of C_M^0 . This is particularly useful, because it allows predictions about the best analytical conditions to be used in order to avoid surface concentration effects with the HMDE. A similar equation is available for mercury film electrodes [5]. Because in many cases r_0 , $\bar{\delta}$, D_{Ox} and D_R do not vary over a

large range, average values of these parameters ($r_0 = 0.04$ cm, $\bar{\delta} = 2 \times 10^{-3}$ cm, $D_{\text{Ox}} = D_{\text{R}} = 10^{-5}$ cm²/s, numerical constant = 0.5) give a useful and widely applicable relationship to calculate the order of magnitude of $C_{\text{M}}^{\text{O}}/C_{\text{M}}$: $C_{\text{M}}^{\text{O}}/C_{\text{M}} \approx 0.25 t_{\text{d}}(\text{s})$. As discussed above, this makes it possible to calculate that in most complexing media, the ligand/metal ratio in the bulk solution must be larger than 1000 in order to avoid a surface effect.

Figure 6 shows that when this condition is not met, the currents measured are at least 30% below those that should be obtained in the absence of a surface effect. An even larger discrepancy is obtained when the diffusion coefficient of the complex is lower than that of the free metal ion. This may be important for determinations of total metal concentration when the ligands of the test solution have not been eliminated previously. It may also be very important in attempts to measure the complexing properties of a metal ion in an unknown medium by stripping polarography. Because the importance of this effect will vary along the stripping polarographic wave, it may be expected that not only the limiting current, but also $E_{1/2}$ and the Tafel slope could be modified.

Influence on measurements of complex stability constants

From Fig. 8, it is possible to estimate the error incurred in $\log \alpha$, when the peak potential, E_{p} , is used to compute α from Eqn. 3 without extrapolating its value to $t_{\text{d}} = 0$. The value for $\Delta \log \alpha = (nF/RT)(\Delta E_{\text{p}}(t = 0) - \Delta E_{\text{p}})$ is given in Fig. 8 for curve 3. For example, it can be seen that for $t_{\text{d}} = 85$ s, an error of 0.1 log unit is made for $\log \alpha$. From all the cases studied here, it can be generalized that $C_{\text{L}}/C_{\text{M}}$ must exceed 1000 if an error smaller than 0.05 log unit is needed. This work suggests that two methods can be used to obtain correct α values, despite the surface effect. First, α can be measured from the shift in E_{p} extrapolated to $t_{\text{d}} = 0$. In cases when $D_{\text{ML}} \neq D_{\text{M}}$, the values of i_{p}^{O} and i_{p}^{c} must also be used in Eqn. 3. Because i_{p}^{O} is also affected by surface effects, $i_{\text{p}}^{\text{O}}/t_{\text{d}}$ must be extrapolated to $t_{\text{d}} = 0$, to give $(i_{\text{p}}^{\text{O}}/t_{\text{d}})_{t=0}$. The ratio $(i_{\text{p}}^{\text{O}}/t_{\text{d}})_{t=0}/(i_{\text{p}}^{\text{c}}/t_{\text{d}})$ must then be introduced into Eqn. 3. The second possibility is to use stripping polarography with medium exchange between deposition and stripping. This procedure has the advantage of being potentially applicable not only to labile complexes, as in the preceding case, but to any kind of complexation reaction, because the peak current is then directly proportional to the quantity of metal deposited during the reduction step. The method is, however, much more delicate to apply with the HMDE, especially in preventing penetration of oxygen, and accounting for any reduction during the medium exchange.

Application to measurements of complexation capacity

Complexation capacity is defined, for natural waters, as the overall concentration of all the complexing sites of the medium able to bind a given metal ion M. It is often obtained by making standard additions of the metal ion to the test medium and measuring the a.s.v. or d.p.a.s.v. peak currents as

a function of C_M . Ideally, before the saturation of the ligands, the increase of i_p is smaller than with standard solutions, whereas a linear increase, with the calibration slope, is observed for an excess of metal ion [24]. The value of C_M at the break corresponds to the complexation capacity. This assumes that the complexes formed are electrochemically inert. For labile complexes, the surface concentration effect can simulate, but must not be confused with, a complexation capacity curve [6]. Indeed, because C_M and t_d play an equivalent role in Eqn. 4, the curve $i_p = f(C_M)$ at constant t_d has a similar shape to the $i_p = f(t_d)$ curve at constant C_M (Fig. 6). Such i_p vs. C_M curves are shown for the Pb/TETA system in Fig. 11, as well as the corresponding i_p vs. t_d curves. The equivalence of t_d and C_M is clearly shown here: for a given value of C_L , the change in i_p depends on $t_d C_M$ only if the other experimental conditions remain constant, i.e., C_M^0 can be changed by variations in both the deposition time and in the value of C_M . A similar curve was reported by Bhat and Weber [13] for the addition of cadmium to fulvic acids, which could be due to the above surface effect. For conditions when $D_{ML} < D_M$, the "break" in the capacity curve (at $t_d = 10$ min in Fig. 6) is even better defined, as i_p decreases in the first part of the curve. This break is related to saturation of the ligand at the electrode surface, but not in the bulk solution. When such a break is interpreted as a bulk complexation capacity, the capacity can be underestimated by a factor of 50–100 [6].

Because it was demonstrated above that t_c is directly proportional to C_L (Eqn. 5), it would be attractive to measure the complexation capacity by titrating the ligand within the diffusion layer by varying t_d , instead of by adding metal ion to the bulk solution. This would have the advantage of keeping the value of C_L/C_M constant (and preferably low) in the bulk solu-

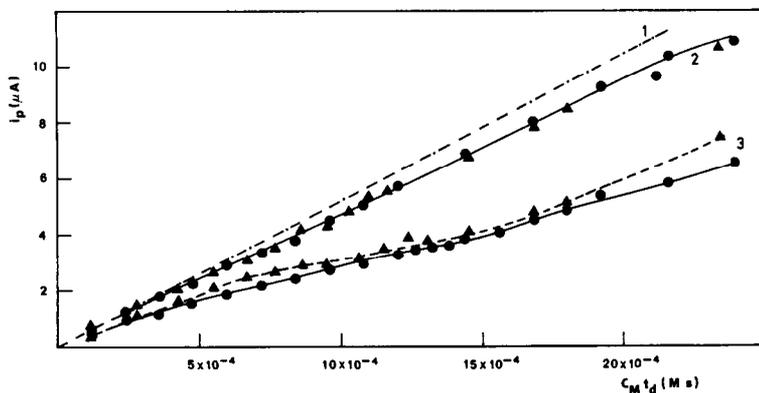


Fig. 11. Change in i_p with t_d and C_M . Curves: (1) $C_L = 0$, pH 5.5 (the same curve was obtained for $C_L \rightarrow \infty$ and pH 8.0); (2) $C_L = 2 \times 10^{-3}$ M and 3×10^{-2} M triethanolamine at pH 8.00; (3) $C_L = 4 \times 10^{-4}$ M and 4×10^{-3} M triethanolamine at pH 8.00. (●) $t_d = 2$ min with variable C_{Pb} ; (▲) $C_{Pb} = 2 \times 10^{-6}$ M with variable t_d . In all cases $E_d = -0.750$ and $v = 20$ mV s $^{-1}$.

tion during the whole procedure, hence minimizing possible coagulation or aggregation effects. Despite its apparent elegance, it is doubtful that such a procedure would give accurate results under practical conditions. One reason is that the measurement of t_c requires the existence of a sufficiently well-defined shoulder on the a.s.v. peak. This occurs in the case of a very stable "simple" complex like Pb(II)/TETA, but not with the mixtures of comparatively weak complexes often found in natural waters. Another operational parameter could be used, such as the jump in the E_p vs. t_d curve, but these "characteristic" parameters would be difficult to find in practical media containing mixtures of ligands with different complexation and diffusing properties.

Surface concentration effect in d.p.a.s.v.

Differential-pulse a.s.v. is more generally used than a.s.v. for application to natural waters, because of its greater sensitivity and its relatively symmetrical peak shapes. Although the modulation in potential during the stripping step will produce a variation of the corresponding concentration of metal ion at the surface for the same t_d , the order of magnitude of the average value of C_M^0 in d.p.a.s.v. will be the same as for a.s.v. However, because of its much greater sensitivity, a smaller t_d can be used with d.p.a.s.v., thus minimizing the surface concentration effect. For practical reasons, it is difficult to use $t_d < 60$ s (i.e., $C_M^0/C_M < 20$) whatever the technique, at least with the HMDE. Thus the surface effect is not negligible, even with d.p.a.s.v., and it can be expected to play an important role in any system containing natural ligands.

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APPENDIX

During the stripping step, the ligand at the electrode surface becomes more and more saturated as the applied potential, E , is scanned to more positive values. This influences the concentration gradients, and in turn, the current and the position of the peak potential. Nevertheless, for $t_d < t_c$ (small decrease in free ligand concentration at the surface compared to the bulk solution, Eqn. 3 can be used as a first approximation by replacing α , with its value at the electrode surface, α_s^0 . $\alpha_s^0 = 1 + \beta[L]^0$, where $[L]^0$ is the concentration of free L at the surface. Further, neglecting the second term in Eqn. 3 and because $\beta[L]^0 \gg 1$, gives $\Delta E_p \approx (RT/nF) \ln(\beta[L]^0)$. $[L]^0$ must be estimated from the diffusion fluxes of L, its protonated forms and ML at the surface. However, if all these species have similar diffusion coefficients, the approximate mass balance equation at the surface will be $C_L^0 = C_L = [L]^0 \alpha_H^0 + [ML]^0$, and $C_M^0 = [M]^0 + [ML]^0$. Because the system is well buffered, $\alpha_H^0 \approx \alpha_H$. Under the conditions used here, L is not saturated and the ML complexes are very stable, so that $C_M^0 \approx [ML]^0$. Finally, C_M^0 is evaluated from Eqn. 4, where the numerical constant is replaced by K' . Combining these conditions with the above equations for ΔE_p , C_L^0 and C_M^0 yields

$$\Delta E_p = (RT/nF) \{ \ln (\beta C_L / \alpha_H) + \ln [1 - (1/K')(D_{Ox}/D_R)^{1/2} (C_M/C_L) t_d / r_0 \bar{\delta}] \} \quad (A1)$$

For $t_d \ll [r_0 \bar{\delta} / (D_{Ox} D_R)^{1/2}] (C_L / C_M)$. Eqn. A1 reduces to a linear relationship

$$\Delta E_p = (RT/nF) \{ \ln (\beta C_L / \alpha_H) - (1/K')(RT/nF)(D_{Ox} D_R)^{1/2} (C_M/C_L) t_d / r_0 \bar{\delta} \} \quad (A2)$$

The ΔE_p values were computed from both Eqns. A1 and A2 with the values of K' , r_0 , $\bar{\delta}$, D_{Ox} and D_R found in this work, and using values of $C_L/C_M = 50, 100$ and 300 . It was found that for the difference in the two ΔE_p values to be less than 2 mV (i.e., to achieve a linear dependence of E_p vs. t_d , the t_d values had to be less than $1, 2$ and 6 min, respectively. In practically similar conditions, the E_p vs. t_d curves were found to be linear over a larger range of t_d , a result that may well be due to the approximate nature of the above mathematical treatment.

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